# SYNTHESIS AND ELECTRICAL PROPERTIES OF THIN FILMS OF FULLERITES ON THE VARIOUS SUBSTRATES

J. Vincenc Oboňa<sup>*a*,1</sup>, Š. Chromik<sup>*a*</sup>, D. Machajdík<sup>*a*</sup>, I. Kostič<sup>*b*</sup>, M. Kadlečíková<sup>*c*</sup>

 <sup>a</sup> Institute of Electrical Engineering, SAS, Dúbravská cesta 9, 841 04 Bratislava, Slovakia
 <sup>b</sup> Institute of Informatics, SAS, Dúbravská cesta 9, 845 07 Bratislava, Slovakia
 <sup>c</sup> Department of Microelectronics FEI STU, Ilkovičova 3, 812 19 Bratislava, Slovakia

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In the process of intercalation,  $C_{60}$  samples in the form of polycrystalline films prepared on various substrates are used. To obtain a homogeneous doping we test two geometrical set-ups of our apparatus, horizontal and vertical. In the present work we focus on the study of transport properties of the samples. We demonstrate the temperature dependence of the conductivity of doped  $C_{60}$  films during the doping process. The time evolutions of the sample resistivity after its exposure to the atmosphere are studied. The structural properties of the samples are examined by X-ray diffraction technique in Bragg-Brentano geometry and Raman Spectroscopy. Scanning Electron Microscopy is used to compare the surface morphology of the undoped and doped  $C_{60}$  films.

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# 1 Introduction

The discovery of the conductivity and the superconductivity [1, 2] in the  $K_x C_{60}$  system opened a question of a potential application of these interesting materials for cryoelectronics. Many authors have focused on the preparation and study of superconducting and transport properties of the  $K_x C_{60}$  system in the ultra-high vacuum at the level of  $10^{-7}$  Pa in situ [1–3], however the physical properties of the sample investigated after breaking the vacuum are not sufficiently known [4].

The  $K_x C_{60}$  reacts very fast on the air. This fact seems to be a serious restriction for the application of the  $K_x C_{60}$ . Because a detailed information about the preparation of the  $K_x C_{60}$  films in the literature is missing or is unclear, our effort was to prepare the  $K_x C_{60}$  films exhibiting a low resistivity. The K (potassium) was intercalated to the  $C_{60}$  film in two different procedures.

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<sup>&</sup>lt;sup>1</sup>E-mail address: Jozef.Vincenc-Obona@savba.sk

We present the resistivity measurements of the  $K_x C_{60}$  as a function of doping and temperature. Some of the prepared films were in situ covered with insulating layers, then exposed to the atmosphere, and the time evolution of the resistivity was studied. We present a possible model of this reaction which explain the changes of the resistivity of the samples.

### 2 Experimental

The undoped C<sub>60</sub> films with a thickness of 70–500 nm were grown by a sublimation of a commercially produced C<sub>60</sub> of 98 % purity in the vacuum. The amount of the deposited material was monitored with a quartz monitor. Ni-Cr pads were evaporated at the corners of the mica (00 $\ell$ ) and Si (111) substrates of the 0.8 × 1.0 cm<sup>2</sup> size. Subsequently, Cu contact wires were attached to the pads with In and covered with silver epoxy. The substrate temperature was measured by Cu-Constantan thermocouple.

To obtain a homogeneous doping we test two geometrical set-ups of our apparatus. In the first case the sample and the source of K are put into a horizontal pyrex tube. Two horizontal ovens ensure the temperature gradient. The  $C_{60}$  film located in the pyrex tube is held at about 200°C and the source of the K is cycled between the room temperature and 150–200°C. To achieve a homogeneous doping (minimum of the resistivity) the films are kept at 200°C for 5–7 hours after each cycle. We used diffblock AV 63 to hold vacuum at the level of  $10^{-3}$  Pa.

In the second case the source of the K and the sample are in the vertical arrangement. The substrates are fixed on the heater and mounted by the contact wire on vacuum feedthrough. The fullerene is evaporated from a Ta boat, the K from a Mo boat, respectively. The distance between the boats and the heater is approximately 10 cm. For the vertical doping the vacuum in the chamber is in the range of  $10^{-4}$  Pa. The substrate with the C<sub>60</sub> thin film is heated up to about  $150^{\circ}$ C and subsequently the K is evaporated on it. This step we repeat several times. When the minimum of the resistivity is achieved, in some cases the insulating layer Na<sub>3</sub>AlF<sub>6</sub>+C<sub>60</sub> is evaporated on the sample.

The structural properties of the samples are examined by X-ray diffraction technique in Bragg-Brentano geometry and Raman Spectroscopy. Scanning Electron Microscopy is used to compare the surface morphology of the undoped and doped  $C_{60}$  films.

### 3 Results and discussion

The fullerene thin film with a tickness of 70–500 nm were evaporated on mica (00 $\ell$ ) and Si (111) substrates. In the process of the sample preparation several parameters, such as the thickness of the C<sub>60</sub> layer, the rate of the deposition and the temperature of the substrates were varied. In all cases the prepared films exhibited a polycrystalline nature.

In the solid state the  $C_{60}$  molecules have a face-centered cubic (fcc) structure and the K is doped into the octahedral or tetrahedral sites [5]. X-ray diffraction patterns show a polycrystalline character of the  $C_{60}$  thin films grown on the mica ( $00\ell$ ) and Si (111) substrates with the preferred (111) and (600) orientation perpendicular to the substrate surface (Fig. 1).

Figure 2 shows SEM images of the mica  $(00\ell)$  surface (Fig. 2a), undoped C<sub>60</sub> film (Fig. 2b) and doped C<sub>60</sub> film (Fig. 2c) on the mica substrate. The layered structure of the mica surface can



Fig. 1. Schematic view of a mica layer, which consists of a hexagonal array of SiO<sub>4</sub> tetrahedra, upon which a C<sub>60</sub> monolayer is then deposited (a) [6]. X-ray diffraction patterns of the 70 nm thick C<sub>60</sub> film grown on the mica ( $00\ell$ ) (b) and on the oxidized Si (111) substrates (c).



Fig. 2. Scanning Electron Microscopy image of the surface of the mica (00 $\ell$ ) substrate (a), the mica (00 $\ell$ )+70 nm C<sub>60</sub> (b) and the mica (00 $\ell$ )+110 nm C<sub>60</sub>+intercalation (c).



Fig. 3. Raman spectra of the 500 nm thick  $C_{60}$  film on the mica (00 $\ell$ ) substrate for three different frequency ranges. Hg (1) mode at 273.5 cm<sup>-1</sup> and Ag (2) mode at 1470.5 cm<sup>-1</sup> for  $C_{60}$  was appeared.

be clearly seen. The polycrystalline  $C_{60}$  film grown on the mica substrate consists of grains with a mean size of approximately 100 nm. After doping the grain coalescence is becoming (Fig. 2c).

We used Raman spectroscopy (Fig. 3) to confirm the existence and to characterise the  $C_{60}$  films. When the film thickness increased up to 500 nm two strong vibrational modes were appeared at 273.5 cm<sup>-1</sup> Hg (1) and at 1470.5 cm<sup>-1</sup> Ag (2) and it is in a good agreement with the value from the literature [7].

In our experiments the pressure was essentially higher at the level of  $10^{-3}$ – $10^{-4}$  Pa comparing to other groups [1–3]. Although we used vacuum background about 5 orders worst, we prepared the samples with the comparable resistance.

We measured the resistivity as a function of exposure time by the two-probe geometry during the doping process. The minimum of the  $\rho$  (exposure time) corresponds to the K<sub>3</sub>C<sub>60</sub> system [8]. It could be appeared the metallic resistivity either in this point or close to this point.

Our samples in the minimum of the resistivity  $(3.23 \times 10^{-3} \,\Omega \text{cm})$  for the vertical arrangement of the experiment have a semiconductor-like temperature dependence because  $K_3C_{60}$  cores of the grains are surrounded by insulating  $K_4C_{60}$  and  $K_6C_{60}$  grain boundaries [8] (Fig. 4). The arrows denote the changes for a fixed temperature increment. The value of the  $\rho$  (exposure time)



Fig. 4. Exposure time dependence of the resistivity of the  $K_x C_{60}$  film during the evaporation of the K at ambient temperature near 25 °C. The  $\rho(0)$  was higher than 320  $\Omega$ cm. The arrows show changes in the resistivity of the sample as it was heated from room temperature to 50 °C.



Fig. 5. The time dependence of the resistivity of the sample after breaking the vacuum.

dependence obtained by four-probe geometry can be different due to contact resistance.

In the case of the horizontal arrangement of the experiment we obtained the lowest resistivity value of  $5.8 \times 10^{-2} \Omega$ cm.

After doping process, at which the minimum of the sample resistivity was achieved, we covered the  $K_x C_{60}$  film (in some cases) with  $C_{60}$  layer (200 nm) and subsequently with Na<sub>3</sub>AlF<sub>6</sub> (500 nm) to limit degradation in atmosphere. The 200 nm thick evaporated  $C_{60}$  layer slowly decreased the sample resistivity due to diffusion of non-stoichiometric potassium from the interstitial sites of the  $K_x C_{60}$ . It seems that this effect can be used for the optimisation of the minimum resistivity value in the future.

In spite of the surface passivation after breaking the vacuum and exposure the sample to the ambient atmosphere the rapid reaction and consecutive changes of the resistivity in the  $K_x C_{60}$  system was detected. We do not see any difference between the passivated and not passivated film concerning of the evolution of the resistivity. It can be due to leakiness of the upper film (pinholes, imperfections). Figure 5 shows the typical resistivity changes as the sample is exposured

to atmosphere.

The left side of the curve in the diagram (Fig. 5) represents the interaction (1) of the  $K_x C_{60}$  with  $O_2$  which is characterised by a sharp increase of the resistivity. However, when potassium attracts from air enough humidity, the reaction of the  $K_x C_{60}$  (K in the interstitial sites of fcc  $C_{60}$ ) with gaseous H<sub>2</sub>O from air is supposed to be dominant (2, 3):

$$K(s) + O_2(g) \to KO_2(s), \tag{1}$$

$$2K(s) + 2H_2O(g) \to 2KOH(aq) + H_2(g), \qquad (2)$$

$$4KO_2(s) + 2H_2O(g,l) \to 4KOH(aq) + 3O_2(g).$$
(3)

Then, the conductivity of KOH (aq) shunts the interaction of the  $K_x C_{60}$  with  $O_2$  (right side).

### 4 Conclusion

We prepared  $C_{60}$  films on mica (00 $\ell$ ) and oxidized Si (111) substrates. X-ray patterns confirm polycrystalline nature of the films. We applied an intercalation of potassium into the fullerene films. We attained the minimum of the resistivity of the samples comparable with the literature in spite of the fact that our experiment were made at standard vacuum conditions ( $\sim 10^{-4}$  Pa). The time dependence of the sample resistivity after breaking the vacuum as well as the possible model of the reaction are added. In the future we will be focused on appropriate methods for the preparation of a stable  $K_x C_{60}$  system on ambient air.

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