# ANALYSIS OF ANELASTIC DEFORMATION RESPONSE IN Fe-B AMORPHOUS SYSTEM<sup>1</sup>

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Anclastic creep aftereffect curves measured on  $Fe_{100-x}B_x$  (x = 14, 17, 20) amorphous alloys under isothermal conditions in the temperature range of 458–637 K have been analysed by two independent direct spectrum analysis methods. Obtained spectra consist of three or more characteristic peaks, each of them represents a type of elementary deformation process. The influence of the structural relaxation, the boron content and the temperature on the shape of the spectra is demonstrated and discussed.

## 1. Introduction

Physical and mechanical properties of amorphous metallic materials – metallic glasses are substantially determined by characteristics of their structural "defects". From the mechanical properties point of view these defects can be characterized as soft regions in amorphous structure which can be deformed much more easily than their surroundings. Different definitions of such defects in amorphous structure have been proposed up to now: (i) free volume concept [1]; (ii) stress inhomogeneities at atomic level [2]; (iii) noncoincident atomic sites configuration at intercluster boundaries [3].

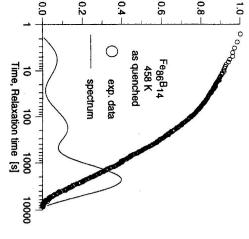
Creep recovery can be considered as the most "pure" anelastic deformation experiment because it takes place without any (or with very low) external stresses. Results and analysis of such experiments seem to be important in crucial appreciation of the models above. It was the reason why we decided to study properties of deformation defects in metallic glasses via the analysis of anelastic deformation response.

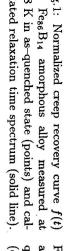
# 2. Experimental and Calculations

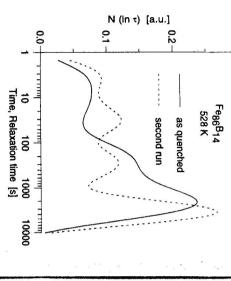
Amorphous ribbons of  $Fe_{100-x}B_x$  (x=14, 17 and 20) were prepared at the Institute of Physics SAS in Bratislava by the rapid quenching process. Samples for creep recovery tests were ground from as received ribbon to the shape usually used for tensile testing with the width of 0.35-0.65 mm and effective length of 12 mm. Measurements were

<sup>&</sup>lt;sup>1</sup>Presented at 9<sup>th</sup> Czech and Slovak conference on magnetism, Košice, Slovakia, August 28-30 1995

0.3







quenched state (solid line) and in second run amorphous alloy measured at 528 K in as-Fig.2: Creep recovery spectra for Fe86 B14 (dashed line).

the second run. rtly relaxed samples the same measurement has been repeated on the same sample noving it creep recovery f(t) was measured. To compare results in as-quenched and about 100 MPa. The stress was kept until steady state creep was reached and on  $58,\,528,\,617$  and 637 K) with the heating rate of 10 K/min and under tensile stress ent resolution of 10 nm. The prepared samples were heated to the test temperature rried out in a SETARAM TMA 92 thermomechanical analyzer which has a displace-

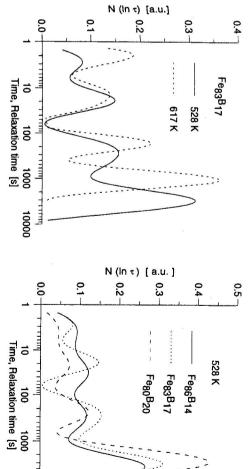
egral torm: netics than a single one, thus in isothermal experiment f(t) can be represented in a Relaxation behaviour in amorphous metallic systems is rather spectrum of simpler

$$f(t) = f_o \int_{-\infty}^{+\infty} N(\ln \tau) \exp(-t/\tau) d\ln \tau, \tag{1}$$

sults. Arrhenius low these methods have been used in our calculations and both of them gave coincident e results to any particular shape of the distribution have been proposed [4, 5]. Both rmalization factor. Up to now two methods for calculation of  $N(\ln \tau)$  without forcing here  $N(\ln au)$  is the relaxation time spectrum, au is the relaxation time and  $f_o$  is the

$$\ln \tau = \ln \tau_o + E/kT,\tag{2}$$

urier transformation [5] ethods. Here we present only results obtained using the method based on discrete d T is the temperature, has been used to compare spectra calculated by these two here  $au_o$  is reverse Debye frequency, E is activation energy, k is the Boltzmann factor



line) and at 617 K (dashed line). amorphous alloy measured at 528 K (solid Fig.3: Creep recovery spectra for Fe<sub>83</sub>B<sub>17</sub>

amorphous alloy measured at 528 K. Fig. 4: Creep recovery spectra for  $Fe_{100-x}B_x$ 

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and for  $\text{Fe}_{40}\text{Ni}_{40}P_{14}\text{B}_{6}$  [10] amorphous alloy. the spectrum is in good agreement with the results obtained for Fe<sub>40</sub>Ni<sub>40</sub>B<sub>20</sub> [6, 7, 8] scale, so they may be a good fit to a lognormal distribution in time. This character of 2400 seconds. Peaks in the relaxation time spectrum are nearly symmetrical on the  $\ln \tau$ finite number of peaks with well defined characteristic relaxation times of 6, 175 and time spectrum computed from experimental data is also shown. Spectrum consists of curve measured at 458 K on Fe<sub>86</sub>B<sub>14</sub> amorphous alloy in as-quenched state. Relaxation Figure 1 shows typical time dependence (in ln t scale) of normalized creep recovery

of two spectra obtained on the same sample in two subsequent runs. It is seen that the with the sharper characteristic relaxation times. relaxation time spectrum in more relaxed state is changed close to the discrete type The same phenomenon we observed in our experiments. Figure 2 shows the comparison has been reported for  $Fe_{40}Ni_{40}B_{20}$  [7] and for  $Ni_{77.5}Si_{17.5}B_{15}$  [9] amorphous ribbons. The decrease of the width of relaxation time peaks during the structural relaxation

shown. Shift of the centers of peaks to the shorter relaxation times is observed for all strated in Fig. 3 where spectra for Fe<sub>83</sub>B<sub>17</sub> measured at two different temperatures are The influence of temperature on the observed relaxation time spectra is demon-

boron concentration in Fig. 4. The character of the spectrum as a superposition of four near lognormal relaxation times ( $\approx 300$  and  $\approx 3000$  s) play more important role in alloys with higher Peaks is preserved for different boron content alloys. Deformation processes with higher The change of the relaxation time spectrum with the chemical composition is shown

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rrangements during deformation and the short-range ordering process in amorphous eak gives new essential information for understanding the mechanisms of atomic re-The fact that the creep recovery spectrum is multimodal rather than a broad single

ident atomic sites configuration at the boundaries of polyclusters. islocation—like atomic defects responsible for plastic deformation. These are noncoinlastic deformation. On the other hand the polycluster model [3] precisely indicates o microscopic notion about dislocations in amorphous structure and accomplish their eformation of amorphous metallic materials exists in the literature [11] but there is Rather good experimental evidences of a dislocation like mechanism of the plastic

a metallic glass ribbon has also a discrete character [12]. finite [3]. We would like to note that the spectrum of positron annihilation lifetimes ypes of atomic rearrangements in the intercluster boundary during some kinetic process pectrum has rather discrete character than a broad one. Namely the number of possible Moreover the polycluster model unlike others can explain why the relaxation time

ach peak represents a distribution of different type of deformation defects. f activation energies) are not observed for all peaks. It shows another evidence that ne case when thermally activated relaxation tends to be controlled by a distribution me and the reduction of the peak width with increasing temperature (so typical for erature. In our experiments the shift of the center of peaks to the shorter relaxation emperature while for the former case the peaks must broaden with decreasing temctivation energy). For the later case the distribution parameter eta is independent on onstant pre-exponential factor) or a distribution of frequency factors (with a constant elaxation tends to be controlled either by a distribution of activation energies (with a endence of peaks width eta one can distinguish (see eq. 2) whether a thermally activated gnormal distribution has been discussed in detail in [13]. From the temperature de-The tendency for relaxation time spectra of thermally activated processes to fit a

narp peak. gns of these stresses in different regions of the ribbon cause a blurring of each relative evidently attributed to the first two mentioned types of internal stresses. Different f the width of peaks in the relaxation time spectra during the structural relaxation ii) local stresses of intercluster boundaries [3] or atomic level stresses [2]. The decrease uenching stresses (acting on the scale of several hundredths of micrometers) [16] and uenching stresses (acting on the scale of the whole sample) [14, 15], (ii) submacroscopic erials obtained by rapid quenching process. This hierarchy consists of: (i) macroscopic There is a hierarchy of internal stresses of different ranges in amorphous metallic ma-

# 5. Conclusion

efects which take part on anelastic deformation process in amorphous structure can loys using methods for relaxation time spectra calculation shows that deformation nalysis of the anelastic creep aftereffect on  $Fe_{100-x}B_x$  (x = 14, 17, 20) amorphous

> characteristic peaks, each of them represents a type of elementary deformation process. tary deformation process. Observed spectra consist of three or more relatively sharp be clearly classified according to the time which they need for realization of an elemen-

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