ELECTRONIC BAND STRUCTURE OF USi31

J.A. Morkowski², A. Szajek

Institute of Molecular Physics, Polish Academy of Sciences Smoluchowskiego 17, 60-179 Poznan, Poland

Received 31 July 1995, accepted 8 February 1996

The electronic band structure of USi₃ is calculated by TB LMTO-ASA method. The 5f electrons form a rather broad band with the center above the Fermi level. The d and p electrons contribute considerably to the density of states at the Fermi level.

1. Introduction

of 5f states. In elemental actinide metals structural and electronic properties can be 5f wave functions of neighbouring actinide atoms in solids leads to the delocalization In contrast to 4f electrons in lanthanides, which in most cases are well localized the 5f system in question. actinide compounds from normal, broad-band non-magnetic metals, to superconductors, 5f-5f overlap and 5f-ligand hybridization decides on a broad variety of behaviours in valence states of the ligand atoms. The interplay between these two main effects, direct 5f electrons is the result of hybridization of the 5f orbitals of actinide atoms with the which 5f electrons are delocalized even if the separation between actinide atoms exceeds delocalization of 5f electrons. However, there exist numerous compounds of actinides in between actinide atoms lower than a certain critical value direct 5f-5f overlap results in In an early systematics known as the Hill rule it was conjectured that for a spacings understood assuming increasing degree of localization with the increasing 5f occupancy. electrons in actinides exhibit a diversified behaviour. The overlap of the more extended knowledge of the band structure is thus essential to characterize and understand a heavy fermion systems and localized moment magnetic systems (see e.g. [1]). The the critical values and there is no direct 5f-5f overlap. In such cases delocalization of

¹Presented at 9th Czech and Slovak conference on magnetism, Košice, Slovakia, August 28-30 1995

²E-mail address: jmorkows@ifmpan.poznan.pl

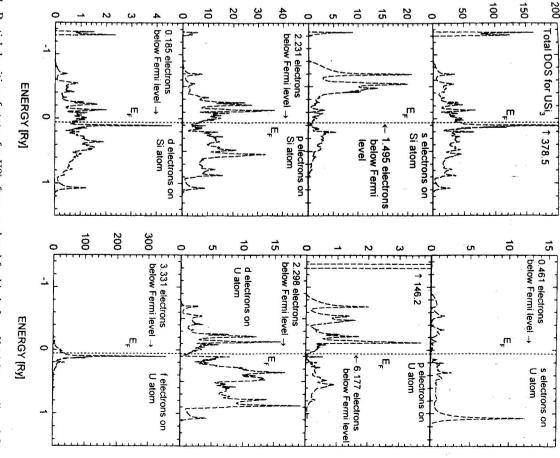
spin-orbit interactions taken into account. The Perdew-Wang [2] potential with nonaverage Wigner-Seitz radius S taken from the experimental value of the lattice constant constant is a=4.0348(8)Å [4]. The band structure calculations were performed for the crystallizes in the Cu₃Au type structure (see e.g. [3] and references therein). The lattice in the Tight Binding Atomic Sphere Approximation (TB LMTO-ASA) [2], with the to compensate for errors due to the atomic spheres approximation. The compound USia The band structure of USi3 was computed using the Linear Muffin Tin Orbital method configurations were taken as: core +5f²6p⁶6d¹7s² for U and core + 3s²3p² for Si 2.8107a.u.giving the overalp volume of muffin-tin spheres 8.9%. The input electronic local corrections was used. The standard [2] combined corrections terms were included the consecutive iteration steps were the same within an error of 0.01 mRy method [5] was employed. The iterations were repeated until the energy eigenvalues of for the fcc system. For integration over the Brillouin zone the modified tetrahedron Computations were done for 286 k-points in the 1/48 wedge of the first Brillouin zone the general condition $4S^3 = 3S_{5i}^3 + S_U^3$. The values were $S_U = 3.3973a.u., S_{5i} =$ the Wigner-Seitz radii S for U and Si atoms were assigned by the program, satisfying $(4\pi/3)S^3 = a^3/4$ for the fcc type unit cell), i.e. for S\(\times 2.9797\) a.u. Different values for

on Si atoms. The high peak appears in the partial density of p states on U, indicating states at the Fermi level is dominated by the f electrons on U atoms and the p-electrons are summarized in the fig.1. The 5f atomic levels of U split into a band of about 0.5 account. No net magnetic polarization and no magnetic moments either on U or Si that p electrons on U occupy essentially the level lying deeply below the Fermi level Ry wide with its gravity center lying about 1eV above the Fermi level. The density of The band structure data calculated for the experimental value of the lattice parameter as a metallic, Pauli type paramagnet atoms appeared within the numerical accuracy of $10^{-4} \mu_B/\text{atom}$. The system behaves We extended the band structure calculations of USi3 with spin polarization taken into

Discussion

is lower here than the value 74 states/(Ry cell) reported in [6] and 65.08 states/(Ry points in the irreducible wedge of the Brillouin zone vs. 84 in [7]) although consistent present results obtained by the TB LMTO method with improved accuracy (for 286-l specific information on the electronic band structure than available in [7] and [8]. The method [6] and by the LMTO method [7]. Recent experimental studies [8] required more the value $(1+\lambda)DOS(E_F) = 80$ states/(Ry cell) determined from the electronic specific f.u.) given in [7]. The high value of the DOS(E_F) can be considered consistent with in details. The density of states at the Fermi level, DOS(E_F)= 44.56 states/(Ry cell) with the general shape of the spectrum with [6] and [7] revealed significant differences The band structure of USi3 was calculated a decade ago by a self-consistent APW

Electronic band structure of USi3



of States [states/(Ry . unit cell)]

atoms. The DOS values for Si are given for three atoms in the cell. Numbers at the arrows electrons (per atom) below the Fermi level is also provided pointing up give heights of the corresponding peaks. The information on the total number of Fig. 1. Partial densities of states for USi3 for s, p, d and f orbitals from U (right panel) and Si

independent measurements, would be small heat [9] provided in the phonon enhancement parameter λ , not known for USi3 from

Acknowledgements. One of us (J.A.M.) thanks Professor R. Troć for a discussion which stimulated the present study. We also thank Drs. B. Nowak and O.J. Zogal for sending the preprint of ref. [8] and for a usefull discussion.

References

- [1] V. Sechovsky, L. Havela: in Ferromagnetic Materials, vol. 4 ed. by E.P. Wohlfarth and K.H.J. Buschow, North Holland, Amsterdam 1988
- [2]H.L. Skriver: The LMTO Method, Springer, New York 1984; J.P. Perdew, J.A. Chevary, S.H. Vosko, K.A. Jackson, M.R. Pederson, D.J. Singh, C. Fiolhais: Phys. Rev. B46 (1992) 6671; G. Krier, O. Jepsen, A. Burkhardt, O.K. Andersen: The TB-LMTO-ASA program (source code version 4.7)
- [3] K. Remschnig, T. Le Bihan, H. Noql, P. Rogl: J. Solid State Chemistry 97 (1992) 391
- [4] H.R. Ott, F. Hulliger, H. Rudigier: Phys. Rev. B31 (1985) 1329
 [5] P.E. Blöchl, O. Jepsen, O.K. Andersen: Phys. Rev. B49 (1994) 16223
- [6] A. Hasegawa: J. Mag. Mag. Mat. 52 (1985) 425
 [7] P. Strange: J. Phys. F: Met. Phys. 16 (1986) 1515
 [8] B. Nowak, O.J. Zogal: Physica B 215 (1995) 193
 [9] M.H. van Maaren, H.J. van Daal, K.H.J. Buschow:
- M.H. van Maaren, H.J. van Daal, K.H.J. Buschow: Solid State Commun. 14 (1974) 145

5