## LETTER TO THE EDITOR

## OBLIQUE INCIDENCE SPECULAR REFLECTIVITY OF UNIAXAL SEMICONDUCTORS NEAR TWO-DIMENSIONAL CRITICAL POINTS

FOLTIN O.1), Bratislava

## КОЭФФИЦИЕНТ ЗЕРКАЛЬНОГО ОТРАЖЕНИЯ ВЕЛИЗИ ДВУМЕРНЫХ КРИТИЧЕСКИХ ТОЧЕК ПРИ НАКЛОННОМ ПАДЕНИИ СВЕТА НА ОДНООСНЫЕ ПОЛУПРОВОДНИКИ

The spectral dependence of specular reflectivity can help in detecting and identifying critical points. So far the cases of normal incidence [1, 2] and oblique incidence [3] reflectivity of isotropic semiconductors and the case of oblique incidence reflectivity of uniaxial semiconductors near three-dimensional critical points [4] have been considered.

Optically uniaxial crystals (composed of the trigonal, tetragonal and hexagonal systems) are characterized by  $\hat{N}_{\perp}$  and  $\hat{N}_{\parallel}$  (complex refractive indices for electromagnetic waves polarized so that the  $\boldsymbol{E}$  vector vibrates normal to and parallel to the optic axis, respectively) or by the respective complex dielectric functions  $\hat{\varepsilon}_{\perp}$  and  $\hat{\varepsilon}_{\parallel}$ . There holds

$$\hat{\mathbf{e}}_{\perp} = \mathbf{e}_{\perp 1} + \mathrm{i}\mathbf{e}_{\perp 2} = \tilde{N}_{\perp}^2 = (n_{\perp} + \mathrm{i}k_{\perp})^2,$$

$$\hat{\epsilon}_1 = \epsilon_1 + i\epsilon_2 = \hat{N}^2 = (n + ik)^2$$
.

(2)

Components of each complex dielectric function (or of each complex refractive index) are related to one another by the Kramers-Kronig dispersion relations. The differences between  $\varepsilon_{1,2}$  and  $\varepsilon_{2}$  arise from differences in the matrix elements governing the transitions. It was shown [4] that the spectral dependence of oblique incidence specular reflectivity for an extraordinary wave reflected by the plane normal to the optic axis (i.e. by the basal plane) can be used for the detection of the singularities of  $\varepsilon_{2}$ , associated with the three-dimensional critical points when  $\varepsilon_{1,2}$  is unsuitable to detect them.

We assume that the critical points are two-dimensional and the spectral dependence of  $\varepsilon_{1,2}$  is unsuitable to detect them. For the direct transitions allowed in the electric dipole approximation we assume the matrix element for  $\varepsilon_2$  to be constant and different from zero and we neglect the weaker  $\omega^{-2}$  factor [5]. Then the frequency dependences of the components of  $\hat{\varepsilon}$  will have near the two-dimensional critical points a step function and logarithmic singularity characters [2] as shown in Fig. 1.

¹) Katedra fyziky, Elektrotechnická fakulta SVŠT, Mlynská dolina, 81219 BRATISLAVA, Czechoslovakia

uniaxial crystal for the extraordinary wave (i.e. E paralell to the plane of incidence) will be [6] In the case of reflection by the basal plane, the oblique incidence specular reflectivity of the

$$R_{\rho} = \frac{(a_{\parallel} - c\cos\varphi)^{2} + (b_{\parallel} - d\cos\varphi)^{2}}{(a_{\parallel} + c\cos\varphi)^{2} + (b_{\parallel} + d\cos\varphi)^{2}},$$

3

where

$$\begin{cases} a_{i}^{2} \\ b_{i}^{2} \end{cases} = [(n_{i}^{2} - k_{i}^{2} - \sin^{2} \varphi)^{2} + 4n_{i}^{2} k_{i}^{2}]^{1/2} \pm (n_{i}^{2} - k_{i}^{2} - \sin^{2} \varphi),$$

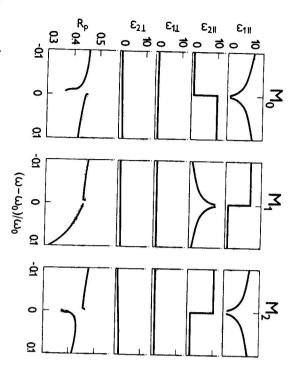
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$$c = n_i n_i - k_i k_i,$$

$$d = n_i k_i + n_i k_i$$

$$d=n_{\perp}k_{\parallel}+n_{\parallel}k_{\perp}$$

and  $\varphi$  is the angle of incidence.



extraordinary wave oblique incidence specular reflectivity  $R_p$ , computed from the frequency depen-Fig. 1. Frequency dependences (near the two-dimensional critical points M<sub>0</sub>, M<sub>1</sub>, M<sub>2</sub>) of the dences of real  $\varepsilon_1$  and imaginary  $\varepsilon_2$  components of the complex dielectric functions  $\hat{\varepsilon}_i$  and  $\hat{\varepsilon}_1$ . The angle of incidence  $\varphi = 80^{\circ}$ .  $\omega_b$  is the frequency of the critical point.

critical points of optically uniaxial crystals. plane reflection (at a large angle of incidence) can be used for the detection of the two-dimensional spectral dependence of the extraordinary wave oblique incidence specular reflectivity at the basal gularities of the components of  $\hat{e}_i$  at critical points give rise to singularities of reflectivity  $R_p$ . So the incidence specular reflectivity  $R_p$  near two-dimensional critical points. Fig. 1 shows that the sinrelations (1) to (6), we computed the spectral dependences of the extraordinary wave oblique nonsingular spectral dependence of the components of  $\hat{e}_1$ , all shown in Fig. 1, and using the Having assumed the above mentioned spectral dependence of the components of  $\hat{e}_{i}$  and the

> correct experimental results. production of such planar surface (normal to the easy-cleavage planes) often prevents to obtain the incidence specular reflectivity in the case of reflection by the plane parallel to the optic axis. But the These critical points could be also detected by means of the spectral dependence of normal

## REFERENCES

- [1] Doran, N. J., Harper, P. G.: phys. stat. sol. (b) 48 (1971), 223.
  [2] Evans, B. L., in: Optical and Electrical Properties. (Physics and Chemistry of Materials with Layered Structures), Vol. 4, Ed. Lee, P. A., D. Reidel Publ. Co., Dordrecht/Boston 1976.
- [3] Foltin, O.: phys. stat. sol. (b) 143 (1987), K151.

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- [4] Foltin, O.: phys. stat. sol. (b) 149 (1988), K73. [5] Harbeke, G., in: Optical Properties of Solids, Ed. Abelès, F., North-Holland, Amsterdam/London 1972.
- [6] Mosteller, L. P. Jr., Wooten, F.: J. Opt. Soc. Amer. 58 (1968), 511

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