

## SPIN-WAVE CONTRIBUTION TO THE HEAT CAPACITY OF NICKEL<sup>1</sup>

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The heat capacity of two high-purity nickel samples has been measured in the range from 1.8 K to 4.2 K. The existence and magnitude of a spin-wave contribution to the heat capacity of nickel has been established in both samples with a higher accuracy than in previously reported measurements.

### ВЛИЯНИЕ СПИНОВЫХ ВОЛН НА ТЕПЛОЕМКОСТЬ НИКЕЛЯ

Измерена теплоемкость двух образцов никеля высокой чистоты в области температур от 1,8 К до 4,2 К. В обоих образцах более точно чем в предыдущих работах определен вклад спиновых волн в теплоемкость никеля.

### 1. INTRODUCTION

The total molar heat capacity of pure nickel can be written theoretically in the form

$$C = \gamma T + \beta T^3 + \alpha T^{5/2} F\left(\frac{T_g}{T}\right), \quad (1)$$

where  $\gamma T$  is the electronic contribution,  $\beta T^3$  is the lattice contribution, and the third term is the spin-wave contribution assuming the presence of spin-waves in ferromagnetic nickel. The coefficient  $\gamma$  is directly proportional to the density of states  $D(\epsilon_F)$  at the Fermi level  $\epsilon_F$

$$\gamma = \frac{\pi^2}{3} k_B^2 D(\epsilon_F),$$

where  $k_B$  is the Boltzmann constant.

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The coefficient  $\beta$  is related to the Debye temperature  $T_D$

$$\beta = \frac{12}{5} \pi^4 R \frac{1}{T_D^3},$$

where  $R$  is the molar gas constant.  
The coefficient  $\alpha$  is given by

$$\alpha = (0.113/z)R(k_B/2J)^{3/2},$$

where  $z = 4$  for nickel and  $2Js$  is a measure of the exchange interaction between neighbouring spins. The function  $F(T_g/T)$  is the correction function which takes into account the energy gap  $k_B T_g$  in the spin-wave spectrum. According to Argyile et al. [1] and Dixon et al. [2]

$$F\left(\frac{T_g}{T}\right) = 1 - 0.39 \frac{T_g}{T} + 0.04 \left(\frac{T_g}{T}\right)^2$$

$T_g = a + cB$ , where for Ni  $a = 0.6$  K,  $c = 1.3$  K T<sup>-1</sup> and  $B$  is the magnetic induction corrected for demagnetization. Dixon et al. [2, 3] have found evidence for a spin-wave contribution to the heat capacity of iron, iron-nickel and copper-nickel alloys. However, their data for a spin-wave contribution to the heat capacity of nickel are ambiguous. Because of the obvious importance to detect the spin-wave contribution to the heat capacity of nickel we decide to study the heat capacity of nickel in the temperature range from 1.8 K to 4.2 K, and also in the magnetic field

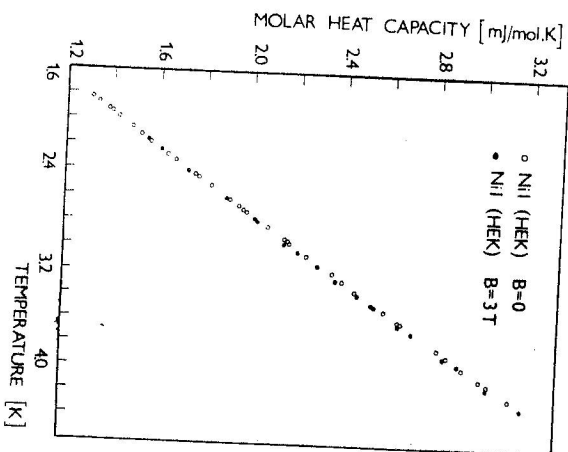


Fig. 1. The temperature and the magnetic field dependence of the molar heat capacity of Ni.

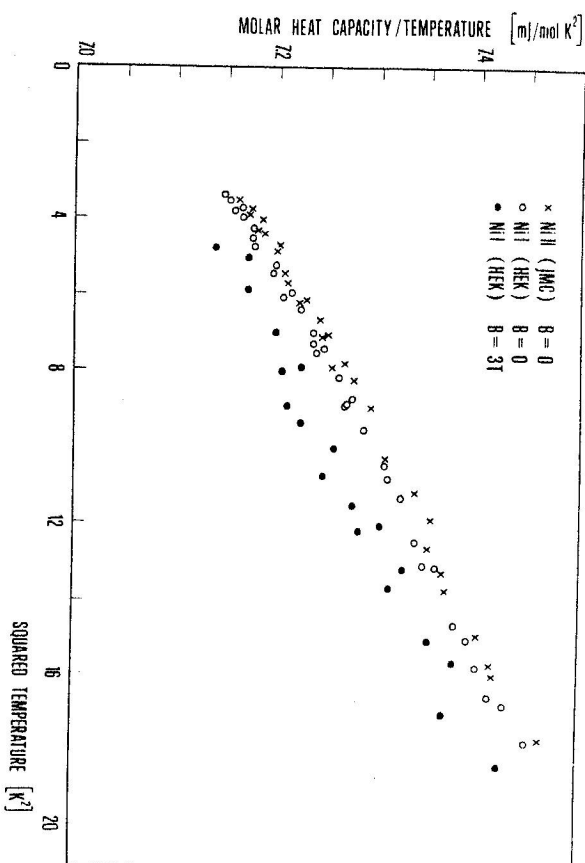


Fig. 2. The dependence of  $C/T$  vs.  $T^2$ .

3 T. The magnetic field quenches the spin-wave heat capacity by introducing a gap in the energy spectrum of the spin-waves.

## II. EXPERIMENTAL DETAILS

The adiabatic pulse method was used. To calibrate the calorimeter, the heat capacity of a pure copper sample was determined. The essential features of the cryostat used in these investigations were published in [4]. The first nickel sample Ni I weighed 18.03352 g and was obtained from HEK, Lübeck, catalogue number M 11 250/64, with the following impurity contents: Fe (2 ppm), Cu (0.5 ppm), Al (1 ppm), Mg (0.5 ppm), Si (3 ppm), Ag (1 ppm). The second sample Ni II weighed 17.29873 g and was obtained from Johnson-Matthey Co., catalogue number JMC 890/8233/1 with the following impurity contents: Fe (10 ppm), Cu (2 ppm), Al (1 ppm), Si (1 ppm), Na (1 ppm), Ca (1 ppm), Ag (1 ppm). Both samples were annealed for 4 hours 40 minutes at 1033 K at a pressure  $10^{-3}$  Pa.

## III. EXPERIMENTAL RESULTS

Fig. 1 shows the data on the temperature dependence of the molar heat capacity of Ni I and the effect of the magnetic field. Three methods of analysis were used to determine the spin-wave contribution.

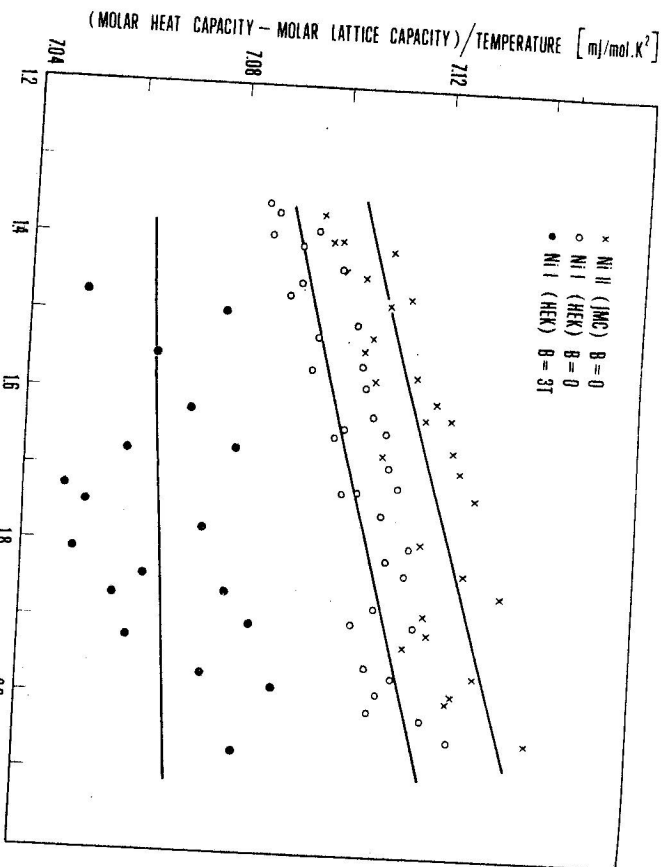


Fig. 3. The dependence of  $(C/T - \beta T^3)$  vs.  $\sqrt{T}$ .

1. The results of measurements were plotted as  $C/T$  vs.  $T^2$  and the effect of a spin-wave term is to cause the experimental curve to be slightly concave towards the  $T^2$  axis because

$$\frac{d^2(C/T)}{d(T^2)^2} < 0$$

for  $T > T_g$ . Fig. 2 shows the data plotted in this way.

2. If we use the recommended value for the Debye temperature  $T_D = 473 \text{ K}$  [5] and the plot  $(C/T - \beta T^3)$  vs.  $\sqrt{T}$ , the slope of such a line with  $T_g = 0$  is the coefficient  $\alpha$ . If  $T_g \neq 0$ , then the slope of such a curve is slightly temperature dependent.

3. In the magnetic field 3 T the spin-waves are quenched and thus the spin-wave contribution is practically eliminated (Fig. 3).

Table 1 shows the result of fitting the experimental data to Eq. (1) with  $T_g = 0.6 \text{ K}$  using the recommended value for the Debye temperature  $T_D = 473 \text{ K}$  [5].

Table 1			
	Ni I		Ni II
$\alpha$ [mJ/mol K <sup>5/2</sup> ]	$0.036 \pm 0.004$		$0.040 \pm 0.006$
$T_D$ [K]	473.0		473.0
$\gamma$ [mJ/mol K <sup>2</sup> ]	$7.050 \pm 0.006$		$7.056 \pm 0.009$

#### IV. CONCLUSION

The experimental results of the heat capacity of two high-purity nickel samples obtained from the different manufactures indicate unambiguously the presence of a spin-wave contribution. The magnitude of the spin-wave contribution in both samples was established with a higher accuracy than in previously reported measurements. The coefficient  $\alpha$  is very close to the value of  $(0.033 \pm 0.004) \text{ mJ/mol K}^{5/2}$  calculated from neutron scattering experiments [6].

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